## Lecture 13: Dynamics of a Particle in a Rotating Coordinate System

$$\vec{F} = m\vec{a} \implies \vec{F}' = m\vec{a}'$$

 $\vec{F}' = \vec{F}_{physical} - m\vec{A}_0 - 2m\vec{w} \times \vec{v}' - m\vec{w} \times (\vec{w} \times \vec{r}') - m\dot{\vec{w}} \times \vec{r}' = m\vec{a}'$ 

- $\vec{A}_0$  is the acceleration of the origin  $\sigma'$  with respect to  $\sigma$ .
- $-2m\vec{w} \times \vec{v}'$  is the coriolis force
- $-m\vec{w} \times (\vec{w} \times \vec{r}')$  is the **centrifugal force**
- $-m\vec{w} \times \vec{r}'$  is the transverse force

The coliolis force, centrifugal force, and traverse force make up the **internal forces**.

Inertial forces for a point moving outward in a merry go round:



## **Rotation of Hurricanes**

North Hemisphere: Clockwise



South Hemisphere: Anticlockwise



## **Example 1**: A rolling wheel



Figure 5.2.6 Rotating coordinates fixed to a rolling wheel.

Acceleration of point P with respect to the ground? Assume the wheel has a radius b.

$$\vec{w} = w\hat{k}' = \frac{v_0}{b}\hat{k}'$$
,  $\begin{cases} \vec{r}' = b\hat{i}' \\ \vec{a}' = 0 \end{cases}$ ,  $\vec{v}' = 0$ 

All terms vanish except centripetal!

Transverse Rule:

$$\vec{a} = \vec{w} \times (\vec{w} \times \vec{r}') = w^2 b \hat{k}' \times (\hat{k}' \times \hat{i}') = -w^2 b \hat{i}' = \left| \frac{v_0^2}{b} (-\hat{i}') \right|$$

**Example 2**: Bicycle with constant speed of radius  $\rho$ . Acceleration of the highest point on one of the wheels with respect to the ground?



We choose a coordinate system with  $\hat{k}'$  always pointing up and  $\hat{i}'$  always along the inner radius of the track. (Note the wheel is rotating with respect to this system).

This system rotates with  $\vec{w} = (v_0/\rho)\hat{k}$ .

 $\begin{cases} \sigma' \text{ does not coincide with } \sigma! \\ \text{Origin } \sigma' \text{ is accelerated: } \vec{A_{\sigma}} = \frac{v_0^2}{\rho} \hat{i}' \end{cases}$ 

The acceleration of any point of the wheel viewed by the  $\sigma'$  system is  $v_0^2/b$  directed at the origin  $\sigma'$ ;

Therefore the top of the wheel:  $\vec{a}_p ' = -(v_0^2/b)\hat{k}'$ Also, the velocity of this point is  $\vec{v}_p ' = -v_0\hat{j}'$ .

So the coriolis acceleration is

$$2\vec{w} \times \vec{v}\,' \ = \ 2\left(\frac{v_0^2}{\rho}\right) \hat{k}\,' \times (-\hat{j}\,') \ = \ \left(\frac{2v_0^2}{\rho}\right) \hat{i}\,'$$

Note:  $\dot{\vec{w}} = 0$  and  $\vec{w} \parallel \vec{r}'$ , so transverse and centripetal are zero! Therefore, the net acceleration relative to the ground of the top of the wheel

$$\vec{a}_p = \vec{a_p'} + 2\vec{w} \times \vec{v}' + \vec{A}_\sigma = \frac{3v_0^2}{\rho}\hat{i}' - \frac{v_0^2}{b}\hat{k}'$$

**Example 3**: Bead on a rod; bead starts at a fixed point with initial velocity  $\vec{v}_0' = w l \hat{i}$ . When does it escape the rod?



We solve in the rotating reference frame:

$$\Sigma F' = N\hat{j}' - 2m\vec{w} \times \vec{v}' - m\vec{w} \times (\vec{w} \times \vec{r}') = m\ddot{\vec{r}}'$$

Note that  $\vec{w} = w\hat{k}'$  and  $\vec{v}' = \dot{x}'\hat{i}'$  and  $\vec{r}' = x'\hat{i}'$  and  $\ddot{\vec{r}}' = \ddot{x}'\hat{i}'$ 

$$N\hat{j}' - 2mw\dot{x}'\hat{j}' - mw^2x'(-\hat{i}') = m\ddot{x}'\hat{i}'$$

 $\implies \begin{cases} N = 2mw\dot{x}' &, \text{ (Coriolis cancels normal force!)} \\ mw^2x' = m\ddot{x}' &, \text{ (Centrifugal force pulls particle to the outside!)} \end{cases}$ 

$$\ddot{x}' = w^2 x'$$

$$x'(t) = Ae^{wt} + Be^{-wt}$$
  
$$\dot{x}'(t) = w(Ae^{wt} - Be^{-wt})$$

$$\begin{aligned} x'(t=0) &= 0 \implies A+B = 0 \implies \boxed{A=-B} \\ \dot{x}'(t=0) &= v'_0 \implies w(A-B) = v_0 ' \implies w(2A) = v_0 ' \implies \boxed{A = \frac{v'_0}{2w}} \quad , \quad \boxed{B = -\frac{v'_0}{2w}} \end{aligned}$$

$$\implies x'(t) = \frac{v'_0}{2w} \left( e^{wt} - e^{-wt} \right) = \frac{v'_0}{w} \sinh(wt)$$

Time to escape rod:

$$x'(t=T) = l = \frac{v'_0}{w}\sinh(wT)$$

$$\implies T = \frac{1}{w}\sinh^{-1}\left(\frac{wl}{v_0'}\right) = \frac{1}{w}\sinh^{-1}(1) = \boxed{\frac{0.88}{w}}$$